Time-of-flight secondary ion mass spectrometry imaging of biological samples with delayed extraction for high mass and high spatial resolutions

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RATIONALE: In Time-of-Flight Secondary Ion Mass Spectrometry (TOF-SIMS), pulsed and focused primary ion beams enable mass spectrometry imaging, a method which is particularly useful to map various small molecules such as lipids at the surface of biological samples. When using TOF-SIMS instruments, the focusing modes of the primary ion beam delivered by liquid metal ion guns can provide either a mass resolution of several thousand or a sub-μm lateral resolution, but the combination of both is generally not possible.

METHODS: With a TOF-SIMS setup, a delayed extraction applied to secondary ions has been studied extensively on rat cerebellum sections in order to compensate for the effect of long primary ion bunches.

RESULTS: The use of a delayed extraction has been proven to be an efficient solution leading to unique features, i.e. a mass resolution up to 10000 at m/z 385.4 combined with a lateral resolution of about 400 nm. Simulations of ion trajectories confirm the experimental determination of optimal delayed extraction and allow understanding of the behavior of ions as a function of their mass-to-charge ratio.

CONCLUSIONS: Although the use of a delayed extraction has been well known for many years and is very popular in MALDI, it is much less used in TOF-SIMS. Its full characterization now enables secondary ion images to be recorded in a single run with a submicron spatial resolution and with a mass resolution of several thousand. This improvement is very useful when analyzing lipids on tissue sections, or rare, precious, or very small size samples. © 2015 The Authors. Rapid Communications in Mass Spectrometry published by John Wiley & Sons Ltd.
(BA) is a mode which produces a very narrow beam diameter, well below one micron, using a pair of electrostatic lenses. The disadvantage of this method is the time width of the primary ion pulses, several tens of nanoseconds on the sample plane, which typically results in unit mass resolution. Another disadvantage is the low current of primary ions that necessitates long analysis times. The alternative High Current Bunched (HCBU) mode consists of using three electrostatic lenses and a primary ion buncher system to be produced, thus enabling very short pulses of less than one nanosecond. In this case, high mass resolution is achieved. The drawback is a larger beam diameter resulting in moderate spatial resolution of around a few microns. When analyzing samples such as tissue sections, high mass resolution and sufficient mass accuracy are needed in order to assign ion peaks, such as those of lipid species, with minimum ambiguity. Thus, the HCBU mode is mandatory. If better spatial resolution is needed, some authors have first recorded high-resolution mass spectra using the HCBU mode, before acquiring secondary ion images using the BA mode. Nevertheless this method is not fully satisfactory because, to reconstruct the ion images, the peak selection in the mass spectra is related to the time burst and consequently the mass separation is not sufficient to avoid the selection of quasi-isobaric species. Therefore, an operating mode combining both a high spatial resolution and high mass resolution is needed.

A method of combining these two features is to use a ‘burst’ mode, where the ion beam has a series of short pulses. In this case, long pulses provided by the BA mode are cut into a burst of successive short ion pulses, resulting in both high mass and high spatial resolutions. The disadvantage of this mode is that each of the short primary ion pulses generates its own secondary ion mass spectrum, separated in time by only 25 ns. This mode is acceptable for samples which only contain a few different species or when only low-mass ions are to be analyzed. For real biological samples, which require the analysis of heavier species such as lipids, the mass spectrum becomes too complex. In addition, the authors indicated that the primary ion current is much lower than in BA mode, thus increasing the time needed for data acquisition. Another method which also combines high mass resolution and high spatial resolution for TOF-SIMS instrumentation has recently been published for accurate measurement of isotopic abundances. In that case the primary ions focusing, termed "collimated burst alignment", enables a beam size of \( \sim 100 \text{ nm} \) to be produced with a mass resolution of several thousands. However, this approach also results in very low primary ion currents, well adapted to avoid dead times in secondary ion detection, but it requires long acquisition times, incompatible with the analysis of lipid species on tissue sections.

An entirely different way which has been investigated is to adapt or build instruments other than the usual reflectron-type mass spectrometers. The concept is to use non-pulsed primary ion beams in order to circumvent the disadvantage in having to choose between mass and spatial resolution. An example of such an approach uses a fullereon ion source fitted to a commercial quadrupole time-of-flight (QTOF) instrument. A second new instrument has been designed, in which secondary ions generated by a DC primary ion beam are sampled using a buncher before being mass analyzed by a specially designed TOF analyzer.

An alternative method to maintain the mass resolution with high spatial resolution is to extract secondary ions from the sample surface with a certain delay after the arrival of primary ions. Delayed extraction was first demonstrated by Wiley and McLaren and is used in MALDI. The Wiley and McLaren device uses a simple technique which compensates by means of pulsed extraction for the ion time-of-flight differences induced by distributions of either initial ion velocity or ion position (but not for both simultaneously). In MALDI, spatial and velocity distributions are correlated and it is possible to find a good set of delay times and strength fields to correct for the initial velocity distribution of ions. In TOF-SIMS it can also be used to compensate for the departure time of secondary ions when using long primary ion bunches as in the present case of the BA mode if the ion velocity distribution is not too large. Several authors have mentioned the use of a delayed extraction: in the case of ions produced after ion/surface collisions in depth profiling of Langmuir-Blodgett and of peptides films in the analysis of eukaryotic cells preserved in trehalose, in the analysis of human skin remains in depth profiling of nanostructured layers, in sputtering of silicate glasses, and in the study of topographic and field effects.

Although already used by several authors, the effect of a delayed extraction on the resulting mass spectra has never been fully described, much less so in the case of biological tissue imaging. In the present work, a detailed study of the different primary ion beam focusing modes of a TOF-SIMS instrument has been performed, with extensive comparisons of mass and spatial resolutions on rat cerebellum sections, and with the aim of obtaining with delayed extraction the best compromise between mass resolution, spatial resolution, and acquisition time.

**EXPERIMENTAL**

**Rat cerebellum sections**

Rat cerebellum sections (16 µm thick) were cut at –20 °C in a CM3050-S cryostat (Leica Microsystems SAS, Nanterre, France) and immediately deposited on silicon wafers (2 in. diameter polished silicon wafers; ACM, Villiers-Saint-Frédéric, France). The sections were then dried under vacuum for 10 min before analysis.

**TOF-SIMS**

Experiments were performed using a TOF-SIMS IV mass spectrometer (see Fig. 1; ION-TOF GmbH, Münster, Germany), located at the Institut de Chimie des Substances Naturelles (CNRS, Gif-sur-Yvette, France). This instrument is equipped with a LMIQ which delivers a Bi\(_n\) \(^{q+} \) ion beam. Primary ions impact the sample with a kinetic energy of 25 keV at an incidence angle of 45°. All the experiments presented were carried out with Bi\(_n\) \(^{q+} \). The emitted secondary ions were accelerated to a kinetic energy of 2 keV (2 keV extraction) toward a field-free region and a single-stage reflectron (first-order compensation). Secondary ions were post-accelerated to a kinetic energy of 10 keV before hitting the detector composed of a micro-channel plate, a scintillator and a photomultiplier. Between two successive
primary ion pulses, a low energy (~20 eV) electron flood gun was used to neutralize the sample surface. Pulsed primary ion currents were measured with a Faraday cup located on the grounded sample holder. For HCBU focusing mode, three lenses were used. For the BA focusing mode, only two lenses were used. Typically, the diameter of the beam in HCBU is 2–5 μm and 400 nm in BA mode.\(^{[47]}\)

Protocol for optimization of mass and spatial resolutions using a delayed extraction

To improve the mass resolution, a delayed extraction approach, already existing in the TOF-SIMS setup, can be used. The extraction voltage is set to 2 kV, the rise time to 40 ns and the delay is adjustable between –5 μs and +20 μs in steps of 5 ns. The delayed extraction was tuned by measuring both the spatial resolution (‘knife edge method’) and the mass resolution.

On the surface of a rat cerebellum section, images in positive and negative ion mode with an analyzed area of 100 μm × 100 μm and with a pixel size of 0.78 μm were recorded with a fluence of 2.25 × 10\(^{11}\) ions-cm\(^{-2}\). The mass resolution of lipid ion peaks was measured for each image from a square (50 μm × 50 μm) in the center of the image. Step-by-step, the extraction delay was increased and for each new extraction delay the settings of the mass spectrometer were optimized. This included adjusting the voltages applied to the X and Y analyzer deflection plates, the voltage on the analyzer lens, and the variable drift path (VDP) voltage (Fig. 1). These optimizations ensured that the secondary ions are correctly collected and time focused on the detector (Fig. 1). These optimizations ensured that the secondary ions are correctly collected and time focused on the detector (Fig. 1).

Furthermore, the delayed extraction voltage modifies the beam raster area. Thus, the voltages applied to the so-called X and Y sensitivity plates (not shown) need also to be adjusted. Each measurement was performed in triplicate.

In order to measure the spatial resolution, images with a fluence of 2 × 10\(^{12}\) ions-cm\(^{-2}\) in both positive and negative ion mode with an analyzed area of 100 μm × 100 μm and a pixel size of 0.2 μm for HCBU and 0.1 μm for BA were recorded on the ground 1000 lines per inch grid located on the sample holder. Line scans were performed perpendicular to the ion beam axis. Scanning the primary ion beam along a straight edge object allows the beam diameter to be measured, provided that the measurement step, which is defined by the pixel size, is much smaller.\(^{[49,50]}\)

Mass spectrometry imaging

On a rat cerebellum section, images in both negative and positive ion mode were acquired to evaluate each acquisition mode. The desired result is to distinguish Purkinje cells, which are very small structures present in the tissue.\(^{[26,51–53]}\) An analyzed area of 400 μm × 400 μm and a pixel size of 0.4 μm were used in BA and in BA plus delayed extraction (BA+DE). In the HCBU operating mode, an analyzed area of 500 μm × 500 μm and a pixel size of 1.9 μm were chosen. The images were recorded with fluence of 5 × 10\(^{11}\) ions-cm\(^{-2}\) in HCBU and of 2 × 10\(^{12}\) ions-cm\(^{-2}\) in BA mode. Due to the very low initial kinetic energy distribution of the secondary ions,\(^{[54]}\) the relationship between the time of flight and the square root of the m/z value is always linear over the whole mass range. The mass calibration was always internal, and signals used for initial calibration were those of H\(^+\), H\(_2\)\(^+\), H\(_3\)\(^+\), C\(^+\), CH\(_2\)\(^+\), CH3\(^+\), C\(_2\)H\(_2\)\(^+\) and C\(_2\)H\(_3\)\(^+\) for the positive ion mode in HCBU and BA acquisition and those of C\(^+\), CH\(_2\)\(^+\), CH\(_3\)\(^+\), C\(_2\)H\(_2\)\(^+\), C\(_3\)\(^+\) and C\(_4\)H\(_2\)\(^+\) for the negative polarity in HCBU and BA acquisition mode. As lighter ions (below m/z 20) in BA+DE acquisition mode are poorly collected by the mass spectrometer, the calibration was still internal with these settings but the signals used were those of C\(_6\)H\(_{12}\)PNO\(_2\)\(^+\) (m/z 224.11), C\(_{27}\)H\(_{45}\)\(^+\) (m/z 369.35), C\(_{27}\)H\(_{45}\)O\(^+\) (m/z 385.35) and C\(_{29}\)H\(_{50}\)O\(_2\)\(^+\) (m/z 430.38) for the positive ion mode, and those of C\(_{18}\)H\(_{35}\)O\(_2\)\(^−\) (m/z 255.23), C\(_{18}\)H\(_{35}\)O\(_2\)\(^−\) (m/z 281.25), C\(_{18}\)H\(_{35}\)O\(_2\)\(^−\) (m/z 283.26), C\(_{27}\)H\(_{45}\)O\(^−\) (m/z 385.35) and C\(_{29}\)H\(_{50}\)O\(_2\)\(^−\) (m/z 429.37) in negative ion mode.

Simulation

The ion trajectories and their TOFs were calculated by computer modeling using SIMION 8.1.1.32-2013-05-20 (Scientific Instrument Services Inc., Ringoes, NJ, USA). The generated ions are evenly distributed from 0 to 100 ns in order to simulate the primary ion pulse duration. The ions are generated at a distance of 1.5 mm from the extraction cone, which is the same as in the instrument. The angular emittance of the ions is ±7° from a Gaussian distribution. The extraction voltage is 2 kV and the initial kinetic energies varied from 0.1 to 1.0 eV.

RESULTS AND DISCUSSION

Delayed extraction optimization

The delayed extraction (DE) was tuned by measuring the mass resolution of several secondary ions detected from the surface of a rat cerebellum section. The optimum delay is determined by a compromise between the mass resolution and the total ion count rate, the peak shape being preserved. In the negative ion mode, mass resolutions of 10000 and 9000 are obtained for the palmitate [C\(_{16}\)H\(_{33}\)O\(_2\)]\(^−\) (m/z 255.2) and cholesterol [M+H] (m/z 385.4) ions, respectively (Fig. 2(a)). In positive ion mode, mass resolutions of 10000 for both the cholesterol [M+H+H\(_2\)O]\(^+\) ion (m/z 369.4) and the vitamin E [M]\(^+\) ion (m/z 430.4) were measured (Fig. 2(b)).
Without DE being applied and in both positive and negative ion mode, the peaks at \( m/z \) 255.2, 369.4, 385.4 and 430.4 exhibit a width (FWHM) of \(~100\) ns, which corresponds to the pulse duration of the primary ions (See Supplementary Fig. S-1, Supporting Information). This parameter strongly limits the mass resolution and the use of DE allows this physical limitation to be overcome.

Another important parameter is the peak shape (Fig. 3). This allows for the exclusion of extraction delays for which a higher mass resolution is obtained but with asymmetric peaks.

The drawback of DE is a lower secondary ion transmission. As shown in Figs. 2(c) and 2(d), the ion peak intensities are half of those obtained without DE. This will be discussed further in the following section.

After optimization, the spatial resolution of each mode has been measured according to the knife edge method (see Supplementary Fig. S-2, Supporting Information). The results are summarized in Table 1. A clear decrease in the beam diameter is observed when comparing HCBU and BA modes. However, a similar beam diameter of \(~0.4\) μm was measured in BA mode with or without the DE indicating that the DE is fully compatible with high lateral resolution. It must be also noted that there is no significant modification of the mass resolution caused by using the DE in the HCBU focusing mode (data not shown).

**Simulation**

Simulations of secondary ion trajectories were performed with the SIMION software. Figure 4 shows the normalized intensity as a function of the calculated TOF for ions with an \( m/z \) of 400 (time bin of 2 ns) with different extraction delays (50 ns, 100 ns, 150 ns, and 200 ns). For an extraction delay of 50 ns, the first emitted ions are well grouped in time, unlike the ions emitted after the extraction potential is applied. This means that the DE only groups the ions emitted during the first 50 ns, as the other ions are shifted as a function of the emission time. If the extraction delay corresponds exactly to the duration of ion production, all the ions are grouped together and a high resolution peak is obtained. Without an
accelerating potential during the arrival of the primary ion pulse, the initial axial velocity of the ions leads to a distribution in the acceleration region and subsequently a distribution in TOF. This distribution in time and distance from the sample seems to be compensated for during the TOF. The compensation is related to the energy fixed by the potential corresponding to their position when the extraction potential is applied and the use of the reflectron. For example, for ions emitted normal to the surface at time zero with an initial energy of 0.5 eV, after 100 ns when the extraction voltage is applied: a $m/z$ 1 ion will be ~1 mm from the surface and the potential at this distance is 1333 V, for a $m/z$ 10 ion this distance will be ~300 μm with a potential of 410 V, and for a $m/z$ 400 ion we obtain a distance of ~50 μm from the surface which would correspond to a potential of 66 V. If the extraction delay is longer than the primary pulse time, the dispersion in position induces a larger difference in energy which does not permit good time compensation, and leads to a tail towards longer TOFs.

The optimum values of the DE, as revealed by these simulations, are significantly smaller than those obtained in the experimental measurements. This is because the instrumental delayed extraction times are calculated using time zero as the pulsation of the primary ions (chopper); the primary ion TOF from the chopper to the sample is thus incorporated into the extraction delay time, while the simulations do not incorporate this time. Therefore, the offset of the experimental extraction delay relative to the simulation values is the summation of the primary ion TOF and the rise time of the extraction pulse.

To determine this offset, the extraction delay was increased step-by-step until the peak of the $H^+$/C0 ion started to shift in time. This phenomenon appears between 1095 and 1100 ns (see Supplementary Fig. S-3, Supporting Information). A linear regression was used to confirm the linear relation between the optimum extraction delay and the primary ion pulse duration. The offset value determined with this method is 1137 ± 8 ns (see Fig. 5). The difference of about 40 ns between these two values could correspond to the rise time of the extraction pulse. Using the value of 1137 ns as the offset between the experiment and the simulation, we find that the simulation would give an

<table>
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<tr>
<th>Table 1. Main features of the three setting modes and consequent optimum parameters for data acquisition</th>
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<td><strong>Focusing mode</strong></td>
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<td>Beam size (μm)</td>
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<td>Extraction delay (ns)</td>
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<td>Pulsed primary ion current (pA)</td>
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<td>Primary ion dose (ions·cm⁻²)</td>
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<td>Acquisition time (h)</td>
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*aMass resolutions have been measured in a center square region of interest of 50 μm × 50 μm, in negative mode on the m/z 385.4 ion peak.
*bPrimary ion current has been measured at a frequency of 10 kHz.

Figure 4. Simulated mass spectra of $m/z$ 400 ions with different initial kinetic energy values (0.1, 0.2, 0.5 and 1.0 eV), produced during a burst length of 100 ns and extracted after a delay of 50 ns (black), 100 ns (red), 150 ns (green), and 200 ns (blue), respectively. Vertical arrows indicate, for 200 ns extraction delay, the groups of ions having initial kinetic energies of 0.2, 0.5, and 1.0 eV, respectively.

Figure 5. Optimum secondary ion extraction delay as a function of the pulse duration of primary ions burst. Equation of the calculated linear regression (solid line), with $E_d$, Extraction delay; $P_d$, Pulse duration; and $R$, coefficient of correlation.
optimal delayed extraction time of 1237 ns for a pulse duration of 100 ns. Recalling that the optimum delayed extraction time determined experimentally was 1235 ns, we find excellent agreement between the simulation and the experiment.

Another simulation to investigate the effect of secondary ion initial kinetic energy on mass resolution was performed, where the analyzer acts as a kinetic energy filter. Supplementary Fig. S-4 (Supporting Information) shows the arrival position of secondary ions on the surface of the micro-channel plate (MCP). The ions with an m/z of 400 and a kinetic energy of 0.1 and 0.2 eV are well focused on the center of the MCP. On the contrary, the ions with larger kinetic energies of 0.5 and 1.0 eV have a larger spot size which is displaced from the center. The increased spot size and modified trajectory are a consequence of the higher initial velocity of these species, which have larger energy deficit due to their position when the extraction potential is applied. This phenomenon could be responsible for the observed decrease in total ion counts, although the secondary ion optics were optimized to maximize the total ion count rate.

Finally, the trajectories of ions at different mass-to-charge ratios with same kinetic energy of 0.1 eV were studied in BA +DE mode and with the optimum value of the extraction delay. Supplementary Fig. S-5 (Supporting Information) shows that trajectories of ions with low mass-to-charge ratio are displaced from the center of the MCP. This result is similar to the previous case as the ions with higher initial velocities have larger energy deficits and thus their trajectories exiting the reflectron are displaced relative to lower kinetic energy ions. This result could explain why lighter ions (lower than m/z 20) are poorly detected with the DE as even low kinetic energy species obtain large energy deficits.

These low-mass ions are usually used in TOF-SIMS for the calibration of the mass spectra. Since the DE induces a kinetic energy deficit depending on the mass-to-charge ratio of the ions, the relationship between the measured TOF and the square root of the m/z is then not linear, especially for light ions, and thus the spectra could not be calibrated using these light mass ions. To obtain a sufficient mass accuracy (usually ~20–50 ppm for TOF-SIMS analysis of tissue samples) for ions of interest, it is better to surround...
the peaks of interest with peaks of known masses, and adding an internal standard should be considered in future experiments.

Application to lipid imaging on rat cerebellum sections

Secondary ion images of the surface of rat cerebrellums were recorded for all acquisition modes; the settings are detailed in Table 1. Figure 6 illustrates the differences in mass resolution between HCBU, BA and BA+DE modes in a selected region of the mass spectra, showing sulfatide ion peaks. The isotope peaks for the sulfatide ions cannot be resolved using the BA settings but are well separated in BA+DE mode. Similar results for fatty acids in negative ion mode and cholesterol and vitamin E in positive ion mode are shown in Supplementary Figs. S-5 and S-7 (Supporting Information), respectively.

In order to illustrate the lateral resolution performances of each focusing mode, special attention was paid on the localization of Purkinje cells, which are localized in the gray matter. These cells have an average size of about 20 μm and exhibit a specific lipid composition. Figure 7 shows three-color overlays of ion images obtained for the [C18:0–H] ion (m/z 283.3) in red, the [C16:0–H] ion (m/z 255.2) in green, and the cholesterol [M–H] ion (m/z 385.4) in blue, acquired in the three different modes. In Fig. 7(a) (HCBU), while it is possible to distinguish the different histological areas of the cerebellum, the Purkinje cells, defined by a higher intensity of the [C18:0–H] ion signal in red, are poorly resolved. In the zoom of Fig. 7(a) (shown in Fig. 7(d)), individual pixels are clearly visible. Figures 7(b) and 7(c) show the images recorded using the BA and BA+DE modes, respectively. Under these conditions, the Purkinje cells are well defined and fine details of the tissue structure can be distinguished, due to a sub-μm lateral resolution.

CONCLUSIONS

Although the use of a delayed extraction has been known for many years and is very popular in MALDI, it has been used much less in TOF-SIMS. Its utilization now enables secondary ion images of molecule-related ions to be recorded with submicron spatial resolution and a mass resolution of several thousand in a single run. Simulation shows that the optimal extraction delay needs to be equal to the primary ion pulse duration, and an excellent agreement is found by experiment. However, the use of the BA+DE setting mode requires a larger primary ion dose density in order to obtain ion images with intensity scales comparable with those recorded with the other modes, such as HCBU.

Another way to obtain high mass resolution and spatial resolution together is to acquire HCBU and BA images on the same area, successively. However, switching from one mode to another takes significant time, especially if a large number of acquisitions is needed, and often the mass assignment is not precise enough with the poor mass resolution obtained in BA mode. In such cases the use of delayed extraction is a better choice, particularly when analyzing rare, precious, or very small samples, such as archeological samples or painting cross-sections.

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REFERENCES


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